

Introduction to fractional calculus

(Based on lectures by R. Gorenflo, F. Mainardi and I. Podlubny)

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July 2008

- Historical origins of fractional calculus
- Fractional integral according to Riemann-Liouville
- Caputo fractional derivative
- Riesz-Feller fractional derivative
- Grünwal-Letnikov
- Integral equations
- Relaxation and oscillation equations
- Fractional diffusion equation
- A nonlinear fractional differential equation. Stochastic solution
- Geometrical interpretation of fractional integration

Fractional Calculus was born in 1695



G.F.A. de L'Hôpital
(1661–1704)

What if the
order will be
 $n = \frac{1}{2}$?

It will lead to a
paradox, from which
one day useful
consequences will be
drawn.

$$\frac{d^n f}{dt^n}$$



G.W. Leibniz
(1646–1716)

G. W. Leibniz (1695–1697)

In the letters to J. Wallis and J. Bernulli (in 1697) Leibniz mentioned the possible approach to fractional-order differentiation in that sense, that for non-integer values of n the definition could be the following:

$$\frac{d^n e^{mx}}{dx^n} = m^n e^{mx},$$

L. Euler (1730)

$$\frac{d^n x^m}{dx^n} = m(m-1) \dots (m-n+1)x^{m-n}$$

$$\Gamma(m+1) = m(m-1) \dots (m-n+1) \Gamma(m-n+1)$$

$$\frac{d^n x^m}{dx^n} = \frac{\Gamma(m+1)}{\Gamma(m-n+1)} x^{m-n}.$$

Euler suggested to use this relationship also for negative or non-integer (rational) values of n . Taking $m = 1$ and $n = \frac{1}{2}$, Euler obtained:

$$\frac{d^{1/2} x}{dx^{1/2}} = \sqrt{\frac{4x}{\pi}} \quad \left(= \frac{2}{\sqrt{\pi}} x^{1/2} \right)$$

S. F. Lacroix adopted Euler's derivation for his successful textbook (*Traité du Calcul Différentiel et du Calcul Intégral*, Courcier, Paris, t. 3, 1819; pp. 409–410).

TRAITÉ ÉLÉMENTAIRE
DE
CALCUL DIFFÉRENTIEL
ET DE
CALCUL INTÉGRAL,
PAR S.-F. LACROIX.

SEPTIÈME ÉDITION,
REVUE ET AUGMENTÉE DE NOTES
Par MM. HERMITE et J.-A. SERRET,
MEMBRES DE L'INSTITUT.

J. B. J. Fourier (1820–1822)

The first step to generalization of the notion of differentiation for **arbitrary functions** was done by J. B. J. Fourier (*Théorie Analytique de la Chaleur*, Didot, Paris, 1822; pp. 499–508).

After introducing his famous formula

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(z) dz \int_{-\infty}^{\infty} \cos(px - pz) dp,$$

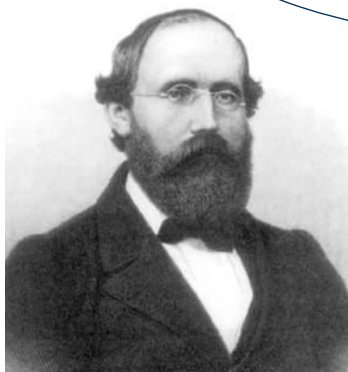
Fourier made a remark that

$$\frac{d^n f(x)}{dx^n} = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(z) dz \int_{-\infty}^{\infty} \cos\left(px - pz + n\frac{\pi}{2}\right) dp,$$

and this relationship could serve as a definition of the n -th order derivative for non-integer n .

Riemann–Liouville definition

$${}_a D_t^\alpha f(t) = \frac{1}{\Gamma(n-\alpha)} \left(\frac{d}{dt} \right)^n \int_a^t \frac{f(\tau) d\tau}{(t-\tau)^{\alpha-n+1}}$$
$$(n-1 \leq \alpha < n)$$



G.F.B. Riemann (1826–1866)



J. Liouville (1809–1882)

Fractional integral according to Riemann-Liouville

- According to Riemann-Liouville the notion of fractional integral of order α ($\alpha > 0$) for a function $f(t)$, is a natural consequence of the well known formula (Cauchy-Dirichlet ?), that reduces the calculation of the n -fold primitive of a function $f(t)$ to a single integral of convolution type

$$J_{a+}^n f(t) := \frac{1}{(n-1)!} \int_a^t (t-\tau)^{n-1} f(\tau) d\tau, \quad n \in \mathbb{N} \quad (1)$$

vanishes at $t = a$ with its derivatives of order $1, 2, \dots, n-1$. Require $f(t)$ and $J_{a+}^n f(t)$ to be *causal* functions, that is, vanishing for $t < 0$.

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- Extend to any positive real value by using the Gamma function, $(n-1)! = \Gamma(n)$

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- Extend to any positive real value by using the Gamma function, $(n-1)! = \Gamma(n)$
- **Fractional Integral of order $\alpha > 0$** (right-sided)

$$J_{a+}^\alpha f(t) := \frac{1}{\Gamma(\alpha)} \int_a^t (t-\tau)^{\alpha-1} f(\tau) d\tau, \quad \alpha \in \mathbb{R} \quad (2)$$

Define $J_{a+}^0 := I$, $J_{a+}^0 f(t) = f(t)$

Fractional integral according to Riemann-Liouville

- Alternatively (left-sided integral)

$$J_{b-}^{\alpha} f(t) := \frac{1}{\Gamma(\alpha)} \int_t^b (\tau - t)^{\alpha-1} f(\tau) d\tau, \quad \alpha \in \mathbb{R}$$

$(a = 0, b = +\infty)$ Riemann $(a = -\infty, b = +\infty)$ Liouville

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- Let

$$J^{\alpha} := J_{0+}^{\alpha}$$

Semigroup properties $J^{\alpha} J^{\beta} = J^{\alpha+\beta}, \quad \alpha, \beta \geq 0$

Commutative property $J^{\beta} J^{\alpha} = J^{\alpha} J^{\beta}$

Effect on power functions

$$J^{\alpha} t^{\gamma} = \frac{\Gamma(\gamma+1)}{\Gamma(\gamma+1+\alpha)} t^{\gamma+\alpha}, \quad \alpha > 0, \gamma > -1, t > 0$$

(Natural generalization of the positive integer properties).

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(Natural generalization of the positive integer properties).

- Introduce the following causal function (vanishing for $t < 0$)

$$\Phi_{\alpha}(t) := \frac{t_+^{\alpha-1}}{\Gamma(\alpha)}, \quad \alpha > 0$$

Fractional integral according to Riemann-Liouville



$$\Phi_{\alpha}(t) * \Phi_{\beta}(t) = \Phi_{\alpha+\beta}(t), \quad \alpha, \beta > 0$$

$$J^{\alpha} f(t) = \Phi_{\alpha}(t) * f(t), \quad \alpha > 0$$

Fractional integral according to Riemann-Liouville



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- Laplace transform

$$\mathcal{L}\{f(t)\} := \int_0^{\infty} e^{-st} f(t) dt = \tilde{f}(s), \quad s \in \mathbb{C}$$

Fractional integral according to Riemann-Liouville



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- Laplace transform

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- Defining the Laplace transform pairs by $f(t) \div \tilde{f}(s)$

$$J^{\alpha} f(t) \div \frac{\tilde{f}(s)}{s^{\alpha}}, \quad \alpha > 0$$

Fractional derivative according to Riemann-Liouville

- Denote by D^n with $n \in \mathbb{N}$, the derivative of order n . Note that

$$D^n J^n = I, \quad J^n D^n \neq I, \quad n \in \mathbb{N}$$

D^n is a left-inverse (not a right-inverse) to J^n . In fact

$$J^n D^n f(t) = f(t) - \sum_{k=0}^{n-1} f^{(k)}(0^+) \frac{t^k}{k!}, \quad t > 0$$

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- Then, define D^α as a left-inverse to J^α . With a positive integer m , $m - 1 < \alpha \leq m$, define:

Fractional Derivative of order α : $D^\alpha f(t) := D^m J^{m-\alpha} f(t)$

$$D^\alpha f(t) := \begin{cases} \frac{d^m}{dt^m} \left[\frac{1}{\Gamma(m-\alpha)} \int_0^t \frac{f(\tau)}{(t-\tau)^{\alpha+1-m}} d\tau \right], & m-1 < \alpha < m \\ \frac{d^m}{dt^m} f(t), & \alpha = m \end{cases}$$

Fractional derivative according to Riemann-Liouville

- Define $D^0 = J^0 = I$.

Then $D^\alpha J^\alpha = I$, $\alpha \geq 0$

$$D^\alpha t^\gamma = \frac{\Gamma(\gamma+1)}{\Gamma(\gamma+1-\alpha)} t^{\gamma-\alpha}, \quad \alpha > 0, \gamma > -1, t > 0$$

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- The fractional derivative $D^\alpha f$ is not zero for the constant function $f(t) \equiv 1$ if $\alpha \notin \mathbb{N}$

$$D^\alpha 1 = \frac{t^{-\alpha}}{\Gamma(1-\alpha)}, \quad \alpha \geq 0, t > 0$$

$1s \equiv 0$ for $\alpha \in \mathbb{N}$, due to the poles of the gamma function

Caputo fractional derivative

- $D_*^\alpha f(t) := J^{m-\alpha} D^m f(t)$ with $m-1 < \alpha \leq m$, namely

$$D_*^\alpha f(t) := \begin{cases} \frac{1}{\Gamma(m-\alpha)} \int_0^t \frac{f^{(m)}(\tau)}{(t-\tau)^{\alpha+1-m}} d\tau, & m-1 < \alpha < m \\ \frac{d^m}{dt^m} f(t), & \alpha = m \end{cases}$$

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- A definition more restrictive than the one before. It requires the absolute integrability of the derivative of order m . In general

$$D^\alpha f(t) := D^m J^{m-\alpha} f(t) \neq J^{m-\alpha} D^m f(t) := D_*^\alpha f(t)$$

unless the function $f(t)$ along with its first $m-1$ derivatives vanishes at $t = 0^+$. In fact, for $m-1 < \alpha < m$ and $t > 0$,

$$D^\alpha f(t) = D_*^\alpha f(t) + \sum_{k=0}^{m-1} \frac{t^{k-\alpha}}{\Gamma(k-\alpha+1)} f^{(k)}(0^+)$$

and therefore, recalling the fractional derivative of the power functions

$$D^\alpha \left(f(t) - \sum_{k=0}^{m-1} \frac{t^k}{k!} f^{(k)}(0^+) \right) = D_*^\alpha f(t), \quad D_*^\alpha 1 \equiv 0, \quad \alpha > 0$$

Riemann versus Caputo



$$D^{\alpha} t^{\alpha-1} \equiv 0, \quad \alpha > 0, t > 0$$

D^{α} is not a right-inverse to J^{α}

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- Functions which for $t > 0$ have the same fractional derivative of order α , with $m-1 < \alpha \leq m$. (the c_j 's are arbitrary constants)

$$D^{\alpha} f(t) = D^{\alpha} g(t) \iff f(t) = g(t) + \sum_{j=1}^m c_j t^{\alpha-j}$$

$$D_{*}^{\alpha} f(t) = D_{*}^{\alpha} g(t) \iff f(t) = g(t) + \sum_{j=1}^m c_j t^{m-j}$$

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- Formal limit as $\alpha \rightarrow (m-1)^+$

$$\alpha \rightarrow (m-1)^+ \implies D^\alpha f(t) \rightarrow D^m J f(t) = D^{m-1} f(t)$$

$$\alpha \rightarrow (m-1)^+ \implies D_*^\alpha f(t) \rightarrow J D^m f(t) = D^{m-1} f(t) - f^{(m-1)}(0^+)$$

- The *Laplace transform*

$$D^\alpha f(t) \div s^\alpha \widetilde{f}(s) - \sum_{k=0}^{m-1} D^k J^{(m-\alpha)} f(0^+) s^{m-1-k}, \quad m-1 < \alpha \leq m$$

Requires the knowledge of the (bounded) initial values of the fractional integral $J^{m-\alpha}$ and of its integer derivatives of order $k = 1, 2, m-1$

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Requires the knowledge of the (bounded) initial values of the fractional integral $J^{m-\alpha}$ and of its integer derivatives of order $k = 1, 2, m-1$

- For the *Caputo fractional derivative*

$$D_*^\alpha f(t) \div s^\alpha \tilde{f}(s) - \sum_{k=0}^{m-1} f^{(k)}(0^+) s^{\alpha-1-k}, \quad m-1 < \alpha \leq m$$

Requires the knowledge of the (bounded) initial values of the function and of its integer derivatives of order $k = 1, 2, m-1$ in analogy with the case when $\alpha = m$

Riesz - Feller fractional derivative

- For functions with Fourier transform

$$\mathcal{F}\{\phi(x)\} = \hat{\phi}(k) = \int_{-\infty}^{\infty} e^{ikx} \phi(x) dx$$

$$\mathcal{F}^{-1}\left\{\hat{\phi}(k)\right\} = \phi(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-ikx} \hat{\phi}(k) dx$$

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- *Symbol* of an operator

$$\hat{A}(k) \hat{\phi}(k) = \int_{-\infty}^{\infty} e^{ikx} A\phi(x) dx$$

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$$\hat{A}(k) \hat{\phi}(k) = \int_{-\infty}^{\infty} e^{ikx} A\phi(x) dx$$

- For the Liouville integral

$$\begin{aligned}J_{\infty+}^{\alpha} f(x) &: = \frac{1}{\Gamma(\alpha)} \int_{-\infty}^x (x - \zeta)^{\alpha-1} f(\zeta) d\zeta \\ J_{\infty-}^{\alpha} f(x) &: = \frac{1}{\Gamma(\alpha)} \int_x^{\infty} (\zeta - x)^{\alpha-1} f(\zeta) d\zeta, \quad \alpha \in \mathbb{R}\end{aligned}$$

Riesz - Feller fractional derivative

- Liouville derivatives ($m - 1 < \alpha < m$)

$$D_{\infty\pm}^{\alpha} = \begin{cases} \pm (D^m J_{\infty\pm}^{m-\alpha}) f(x), & m \text{ odd} \\ (D^m J_{\infty\pm}^{m-\alpha}) f(x), & m \text{ even} \end{cases}$$

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- Operator symbols

$$\hat{J}_{\infty\pm}^{\alpha} = |k|^{-\alpha} e^{\pm i(\text{sign} k)\alpha\pi/2} = (\mp ik)^{-\alpha}$$

$$\hat{D}_{\infty\pm}^{\alpha} = |k|^{+\alpha} e^{\mp i(\text{sign} k)\alpha\pi/2} = (\mp ik)^{+\alpha}$$

$$\hat{J}_{\infty+}^{\alpha} + \hat{J}_{\infty-}^{\alpha} = \frac{2 \cos(\alpha\pi/2)}{|k|^{\alpha}}$$

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$$\hat{J}_{\infty+}^{\alpha} + \hat{J}_{\infty-}^{\alpha} = \frac{2 \cos(\alpha\pi/2)}{|k|^{\alpha}}$$

- Define a symmetrized version

$$I_0^{\alpha} f(x) = \frac{J_{\infty+}^{\alpha} f + J_{\infty-}^{\alpha} f}{2 \cos(\alpha\pi/2)} = \frac{1}{2\Gamma(\alpha) \cos(\alpha\pi/2)} \int_{-\infty}^{\infty} |x - \xi|^{\alpha-1} f(\xi) d\xi$$

(with exclusion of odd integers). The operator symbol

is $\hat{I}_0^{\alpha} = |k|^{-\alpha}$

Riesz-Feller fractional derivative

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- Define the *Riesz fractional derivative* by analytical continuation

$$\mathcal{F}\{D_0^\alpha f\}(k) := \mathcal{F}\{-I_0^{-\alpha} f\}(k) = -|k|^\alpha \hat{f}(k)$$

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generalized by Feller

- D_θ^α = Riesz-Feller fractional derivative of order α and skewness θ

$$\mathcal{F}\{D_\theta^\alpha f\}(k) := -\psi_\alpha^\theta(k) \hat{f}(k)$$

with

$$\psi_\alpha^\theta(k) = |k|^\alpha e^{i(\operatorname{sign} k)\theta\pi/2}, \quad 0 < \alpha \leq 2, |\theta| \leq \min\{\alpha, 2 - \alpha\}$$

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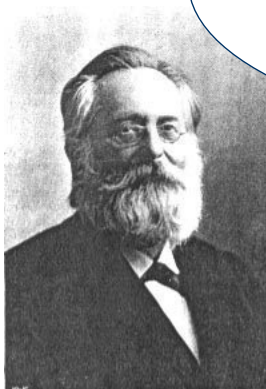
$$\psi_\alpha^\theta(k) = |k|^\alpha e^{i(\operatorname{sign} k)\theta\pi/2}, \quad 0 < \alpha \leq 2, |\theta| \leq \min\{\alpha, 2 - \alpha\}$$

- The symbol $-\psi_\alpha^\theta(k)$ is the logarithm of the characteristic function of a Lévy stable probability distribution with index of stability α and asymmetry parameter θ

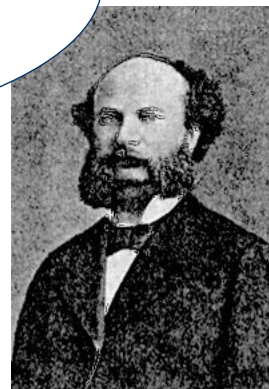
Grünwald–Letnikov definition

$${}_a D_t^\alpha f(t) = \lim_{h \rightarrow 0} h^{-\alpha} \sum_{j=0}^{\left[\frac{t-a}{h} \right]} (-1)^j \binom{\alpha}{j} f(t - jh)$$

$[x]$ – integer part of x



A.K. Grünwald



A.V. Letnikov

- From

$$D\phi(x) = \lim_{h \rightarrow 0} \frac{\phi(x) - \phi(x-h)}{h}$$

...

$$D^n = \lim_{h \rightarrow 0} \frac{1}{h^n} \sum_{k=0}^n (-1)^k \binom{n}{k} \phi(x - kh)$$

- From

$$D\phi(x) = \lim_{h \rightarrow 0} \frac{\phi(x) - \phi(x-h)}{h}$$

$$D^n = \lim_{h \rightarrow 0} \frac{1}{h^n} \sum_{k=0}^n (-1)^k \binom{n}{k} \phi(x - kh)$$

- the Grünwal-Letnikov fractional derivatives are

$${}_{GL}D_{a+}^{\alpha} = \lim_{h \rightarrow 0} \frac{1}{h^{\alpha}} \sum_{k=0}^{[(x-a)/h]} (-1)^k \binom{\alpha}{k} \phi(x - kh)$$

$${}_{GL}D_{b-}^{\alpha} = \lim_{h \rightarrow 0} \frac{1}{h^{\alpha}} \sum_{k=0}^{[(b-x)/h]} (-1)^k \binom{\alpha}{k} \phi(x + kh)$$

$[\bullet]$ denotes the integer part

- **Abel's equation (1st kind)**

$$\frac{1}{\Gamma(\alpha)} \int_0^t \frac{u(\tau)}{(t-\tau)^{1-\alpha}} d\tau = f(t), \quad 0 < \alpha < 1$$

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- The mechanical problem of the *tautochrone*, that is, determining a curve in the vertical plane, such that the time required for a particle to slide down the curve to its lowest point is independent of its initial placement on the curve.

- **Abel's equation (1st kind)**

$$\frac{1}{\Gamma(\alpha)} \int_0^t \frac{u(\tau)}{(t-\tau)^{1-\alpha}} d\tau = f(t), \quad 0 < \alpha < 1$$

- The mechanical problem of the *tautochrone*, that is, determining a curve in the vertical plane, such that the time required for a particle to slide down the curve to its lowest point is independent of its initial placement on the curve.
- Found many applications in diverse fields:
 - Evaluation of spectroscopic measurements of cylindrical gas discharges
 - Study of the solar or a planetary atmosphere
 - Star densities in a globular cluster
 - Inversion of travel times of seismic waves for determination of terrestrial sub-surface structure
 - Inverse boundary value problems in partial differential equations

Abel's equation

- - Heating (or cooling) of a semi-infinite rod by influx (or efflux) of heat across the boundary into (or from) its interior

$$u_t - u_{xx} = 0, \quad u = u(x, t)$$

in the semi-infinite intervals $0 < x < \infty$ and $0 < t < \infty$. Assume initial temperature, $u(x, 0) = 0$ for $0 < x < \infty$ and given influx across the boundary $x = 0$ from $x < 0$ to $x > 0$,

$$-u_x(0, t) = p(t)$$

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- Then,

$$u(x, t) = \frac{1}{\sqrt{\pi}} \int_0^t \frac{p(\tau)}{\sqrt{t-\tau}} e^{-x^2/[4(t-\tau)]} d\tau, \quad x > 0, t > 0$$

Abel's equation (1st kind)



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$$J^\alpha u(t) = f(t)$$

and consequently is solved by

$$u(t) = D^\alpha f(t)$$

using $D^\alpha J^\alpha = I$. Let us now solve using the Laplace transform

$$\frac{\tilde{u}(s)}{s^\alpha} = \tilde{f}(s) \implies \tilde{u}(s) = s^\alpha \tilde{f}(s)$$

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- The solution is obtained by the inverse Laplace transform: Two possibilities :

Abel's equation (1st kind)

- 1)

$$\tilde{u}(s) = s \left(\frac{\tilde{f}(s)}{s^{1-\alpha}} \right)$$

$$u(t) = \frac{1}{\Gamma(1-\alpha)} \frac{d}{dt} \int_0^t \frac{f(\tau)}{(t-\tau)^\alpha} d\tau$$

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- 2)

$$\tilde{u}(s) = \frac{1}{s^{1-\alpha}} [s\tilde{f}(s) - f(0^+)] + \frac{f(0^+)}{s^{1-\alpha}}$$

$$u(t) = \frac{1}{\Gamma(1-\alpha)} \int_0^t \frac{f'(\tau)}{(t-\tau)^\alpha} d\tau + f(0^+) \frac{t^{-\alpha}}{\Gamma(1-\alpha)}$$

Solutions expressed in terms of the fractional derivatives D^α and D_*^α , respectively

Abel's equation (2nd kind)



$$u(t) + \frac{\lambda}{\Gamma(\alpha)} \int_0^t \frac{u(\tau)}{(t-\tau)^{1-\alpha}} d\tau = f(t), \quad \alpha > 0, \lambda \in \mathbb{C}$$

Abel's equation (2nd kind)



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- In terms of the fractional integral operator

$$(1 + \lambda J^\alpha) u(t) = f(t)$$

solved as

$$u(t) = (1 + \lambda J^\alpha)^{-1} f(t) = \left(1 + \sum_{n=1}^{\infty} (-\lambda)^n J^{\alpha n} \right) f(t)$$

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- Noting that

$$J^{\alpha n} f(t) = \Phi_{\alpha n}(t) * f(t) = \frac{t_+^{\alpha n - 1}}{\Gamma(\alpha n)} * f(t)$$

$$u(t) = f(t) + \left(\sum_{n=1}^{\infty} (-\lambda)^n \frac{t_+^{\alpha n - 1}}{\Gamma(\alpha n)} \right) * f(t)$$

Abel's equation (2nd kind)

- Relation to the Mittag-Leffler functions

$$e_{\alpha}(t; \lambda) := E_{\alpha}(-\lambda t^{\alpha}) = \sum_{n=0}^{\infty} \frac{(-\lambda t^{\alpha})^n}{\Gamma(\alpha n + 1)}, \quad t > 0, \alpha > 0, \lambda \in \mathbb{C}$$

$$\sum_{n=1}^{\infty} (-\lambda)^n \frac{t_+^{\alpha n - 1}}{\Gamma(\alpha n)} = \frac{d}{dt} E_{\alpha}(-\lambda t^{\alpha}) = e'_{\alpha}(t; \lambda), \quad t > 0$$

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- Finally,

$$u(t) = f(t) + e'_{\alpha}(t; \lambda)$$

- **Relaxation and oscillation equations. Integer order**

$$u'(t) = -u(t) + q(t)$$

the solution, under the initial condition $u(0^+) = c_0$, is

$$u(t) = c_0 e^{-t} + \int_0^t q(t - \tau) e^{-\tau} d\tau$$

Fractional differential equations

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- For the *oscillation* differential equation

$$u''(t) = -u(t) + q(t)$$

the solution, under the initial conditions $u(0^+) = c_0$ and $u'(0^+) = c_1$, is

$$u(t) = c_0 \cos t + c_1 \sin t + \int_0^t q(t - \tau) \sin \tau d\tau$$

Relaxation and oscillation equations

Fractional version

$$D_*^\alpha u(t) = D^\alpha \left(\left(u(t) - \sum_{k=0}^{m-1} \frac{t^k}{k!} u^{(k)}(0^+) \right) \right) = -u(t) + q(t), \quad t > 0$$

$m-1 < \alpha \leq m$, initial values $u^{(k)}(0^+) = c_k$, $k = 0, \dots, m-1$. When α is the integer m

$$u(t) = \sum_{k=0}^{m-1} c_k u_k(t) + \int_0^t q(t-\tau) u_\delta(\tau) d\tau$$

$u_k(t) = J^k u_0(t)$, $u_k^{(h)}(0^+) = \delta_{kh}$, $h, k = 0, \dots, m-1$, $u_\delta(t) = -u'_0(t)$

The $u_k(t)$'s are the *fundamental solutions*, linearly independent solutions of the *homogeneous* equation satisfying the initial conditions. The function $u_\delta(t)$, which is convoluted with $q(t)$, is the *impulse-response solution* of the *inhomogeneous* equation with $c_k \equiv 0$, $k = 0, \dots, m-1$, $q(t) = \delta(t)$. For ordinary relaxation and oscillation, $u_0(t) = e^{-t} = u_\delta(t)$ and $u_0(t) = \cos t$, $u_1(t) = J u_0(t) = \sin t = \cos(t - \pi/2) = u_\delta(t)$.

Relaxation and oscillation equations

- **Solution of the fractional equation by Laplace transform**

Applying the operator J^α to the fractional equation

$$u(t) = \sum_{k=0}^{m-1} \frac{t^k}{k!} - J^\alpha u(t) + J^\alpha q(t)$$

Laplace transforming yields

$$\tilde{u}(s) = \sum_{k=0}^{m-1} \frac{1}{s^{k+1}} - \frac{1}{s^\alpha} \tilde{u}(s) + \frac{1}{s^\alpha} \tilde{q}(s)$$

hence

$$\tilde{u}(s) = \sum_{k=0}^{m-1} \frac{s^{\alpha-k-1}}{s^\alpha + 1} + \tilde{q}(s)$$

Relaxation and oscillation equations

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- Introducing the Mittag-Leffler type functions

$$e_\alpha(t) \equiv e_\alpha(t; 1) := E_\alpha(-t^\alpha) \div \frac{s^{\alpha-1}}{s^\alpha + 1}$$

Relaxation and oscillation equations



$$u_k(t) := J^k e_\alpha(t) \div \frac{s^{\alpha-k-1}}{s^\alpha + 1}, \quad k = 0, 1, \dots, m-1$$

we find

$$u(t) = \sum_{k=0}^{m-1} u_k(t) - \int_0^t q(t-\tau) u'_0(\tau) d\tau$$

Relaxation and oscillation equations



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- When α is not integer, $m-1$ represents the integer part of α ($[\alpha]$) and m the number of initial conditions necessary and sufficient to ensure the uniqueness of the solution $u(t)$. The m functions $u_k(t) = J^k e_\alpha(t)$ with $k = 0, 1, \dots, m-1$ represent those particular solutions of the *homogeneous* equation which satisfy the initial conditions

$$u_k^{(h)}(0^+) = \delta_{kh}, \quad h, k = 0, 1, \dots, m-1$$

and therefore they represent the *fundamental solutions* of the fractional equation. Furthermore, the function $u_\delta(t) = -e'_\alpha(t)$ represents the *impulse-response solution*.

Fractional diffusion equation

- Fractional diffusion equation, obtained from the standard diffusion equation by replacing the second-order space derivative with a Riesz-Feller derivative of order $\alpha \in (0, 2]$ and skewness θ and the first-order time derivative with a Caputo derivative of order $\beta \in (0, 2]$

$${}_x D_{\theta}^{\alpha} u(x, t) = {}_t D_{*}^{\beta} u(x, t), \quad x \in \mathbb{R}, t \in \mathbb{R}^{+}$$

$$0 < \alpha \leq 2, |\theta| \leq \min\{\alpha, 2 - \alpha\}, \quad 0 < \beta \leq 2$$

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- Space-fractional diffusion $\{0 < \alpha \leq 2, \beta = 1\}$
Time-fractional diffusion $\{\alpha = 2, 0 < \beta \leq 2\}$
Neutral-fractional diffusion $\{0 < \alpha = \beta \leq 2\}$

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Neutral-fractional diffusion $\{0 < \alpha = \beta \leq 2\}$
- Riesz-Feller space-fractional derivative*

$$\mathcal{F} \{ {}_x D_\theta^\alpha f(x); \kappa \} = -\psi_\alpha^\theta(\kappa) \hat{f}(\kappa)$$

$$\psi_\alpha^\theta(\kappa) = |\kappa|^\alpha e^{i(\text{sign } \kappa)\theta\pi/2}, \quad 0 < \alpha \leq 2, |\theta| \leq \min\{\alpha, 2 - \alpha\}$$

Fractional diffusion equation

- *Caputo time-fractional derivative*

$$D_*^\alpha f(t) := \begin{cases} \frac{1}{\Gamma(m-\alpha)} \int_0^t \frac{f^{(m)}(\tau)}{(t-\tau)^{\alpha+1-m}} d\tau, & m-1 < \alpha < m \\ \frac{d^m}{dt^m} f(t), & \alpha = m \end{cases}$$

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- Cauchy problem

$$u(x, 0) = \varphi(x), \quad x \in \mathbb{R}, u(\pm\infty, t) = 0, t > 0$$

$$u_{\alpha,\beta}^\theta(x, t) = \int_{-\infty}^{+\infty} G_{\alpha,\beta}^\theta(\xi, t) \varphi(x - \xi) d\xi$$

$$G_{\alpha,\beta}^\theta(ax, bt) = b^{-\gamma} G_{\alpha,\beta}^\theta(ax/b^\gamma, t), \quad \gamma = \beta/\alpha$$

Fractional diffusion equation

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$$G_{\alpha,\beta}^\theta(ax, bt) = b^{-\gamma} G_{\alpha,\beta}^\theta(ax/b^\gamma, t), \quad \gamma = \beta/\alpha$$

- *Similarity variable x/t^γ*

$$G_{\alpha,\beta}^\theta(x, t) = t^{-\gamma} K_{\alpha,\beta}^\theta(x/t^\gamma), \quad \gamma = \beta/\alpha$$

Fractional diffusion equation

- Solution by Fourier transform for the space variable and the Laplace transform for the time variable

$$-\psi_{\alpha}^{\theta}(\kappa) \widehat{\widehat{G_{\alpha,\beta}^{\theta}}} = s^{\beta} \widehat{\widehat{G_{\alpha,\beta}^{\theta}}} - s^{\beta-1}$$

$$\widehat{\widehat{G_{\alpha,\beta}^{\theta}}} = \frac{s^{\beta-1}}{s^{\beta} + \psi_{\alpha}^{\theta}(\kappa)}$$

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$$\widehat{\widehat{G_{\alpha,\beta}^{\theta}}} = \frac{s^{\beta-1}}{s^{\beta} + \psi_{\alpha}^{\theta}(\kappa)}$$

- Inverse Laplace transform

$$\widehat{\widehat{G_{\alpha,\beta}^{\theta}}}(k, t) = E_{\beta} \left[-\psi_{\alpha}^{\theta}(\kappa) t^{\beta} \right], \quad E_{\beta}(z) := \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(\beta n + 1)}$$

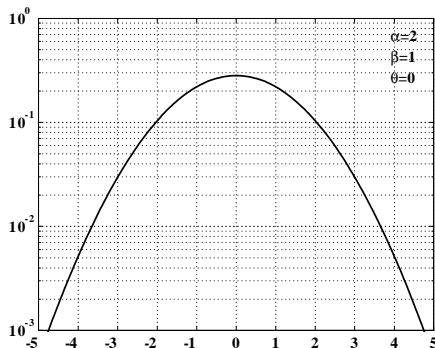
$$G_{\alpha,\beta}^{\theta}(x, t) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{-ikx} E_{\beta} \left[-\psi_{\alpha}^{\theta}(\kappa) t^{\beta} \right] dk$$

Fractional diffusion equation

Particular cases

$\{\alpha = 2, \beta = 1\}$ (*Standard diffusion*)

$$G_{2,1}^0(x, t) = t^{-1/2} \frac{1}{2\sqrt{\pi}} \exp[-x^2/(4t)]$$



Fractional diffusion equation

$\{0 < \alpha \leq 2, \beta = 1\}$ (*Space fractional diffusion*)

The Mittag-Leffler function reduces to the exponential function and we obtain a characteristic function of the class $\{L_\alpha^\theta(x)\}$ of Lévy strictly stable densities

$$\widehat{L}_\alpha^\theta(\kappa) = e^{-\psi_\alpha^\theta(\kappa)}, \quad \widehat{G}_{\alpha,1}^\theta(\kappa, t) = e^{-t\psi_\alpha^\theta(\kappa)}$$

The Green function of the space-fractional diffusion equation can be interpreted as a Lévy strictly stable *pdf*, evolving in time

$$G_{\alpha,1}^\theta(x, t) = t^{-1/\alpha} L_\alpha^\theta(x/t^{1/\alpha}), \quad -\infty < x < +\infty, t \geq 0$$

Particular cases:

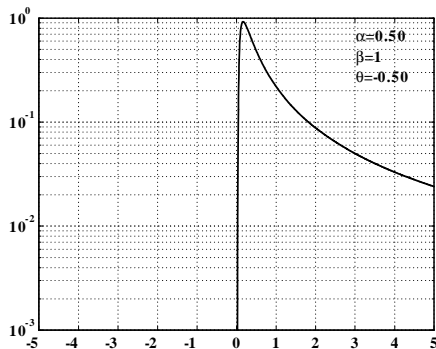
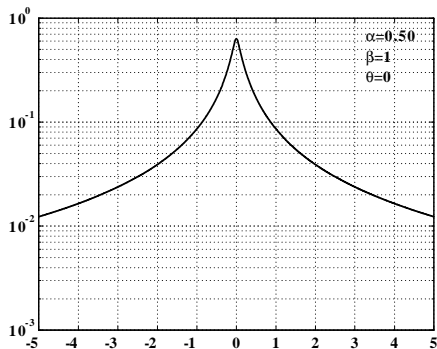
$\alpha = 1/2, \theta = -1/2$, Lévy-Smirnov

$$e^{-s^{1/2}} \stackrel{\mathcal{L}}{\longleftrightarrow} L_{1/2}^{-1/2}(x) = \frac{x^{-3/2}}{2\sqrt{\pi}} e^{-1/(4x)}, \quad x \geq 0$$

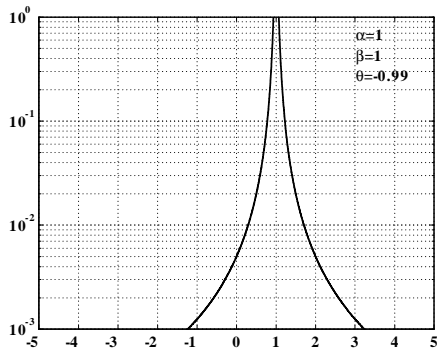
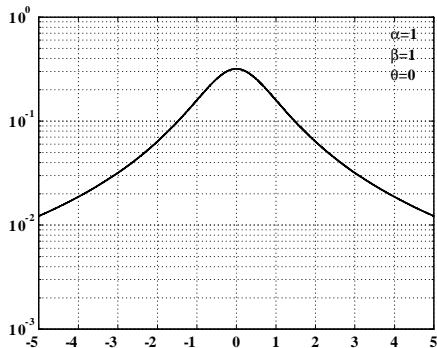
$\alpha = 1, \theta = 0$, Cauchy

$$e^{-|\kappa|} \stackrel{\mathcal{F}}{\longleftrightarrow} L_1^0(x) = \frac{1}{\pi} \frac{1}{x^2 + 1}, \quad -\infty < x < +\infty$$

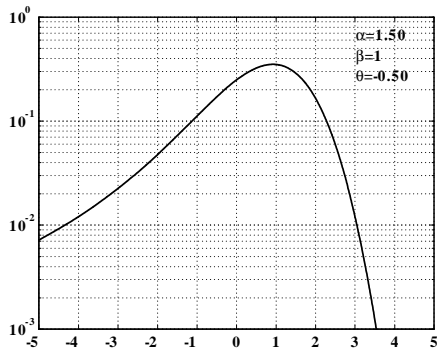
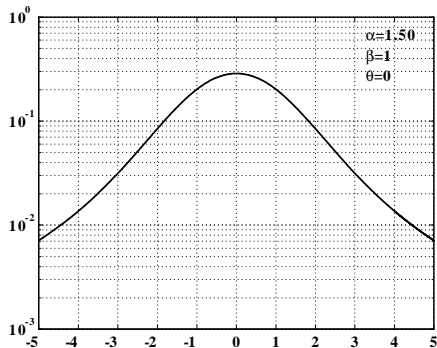
Fractional diffusion equation



Fractional diffusion equation



Fractional diffusion equation



Fractional diffusion equation

$\{\alpha = 2, 0 < \beta < 2\}$ (*Time-fractional diffusion*)

$$\widehat{G_{2,\beta}^0}(\kappa, t) = E_\beta \left(-\kappa^2 t^\beta \right), \quad \kappa \in \mathbb{R}, t \geq 0$$

or with the equivalent Laplace transform

$$\widetilde{G_{2,\beta}^0}(x, s) = \frac{1}{2} s^{\beta/2-1} e^{-|x|s^{\beta/2}}, \quad -\infty < x < +\infty, \Re(s) > 0$$

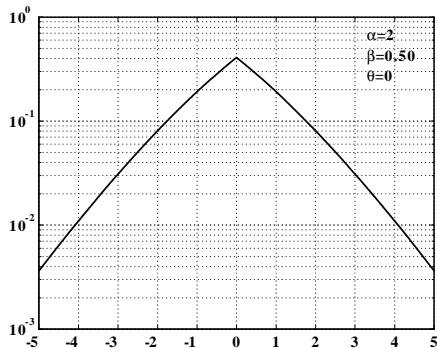
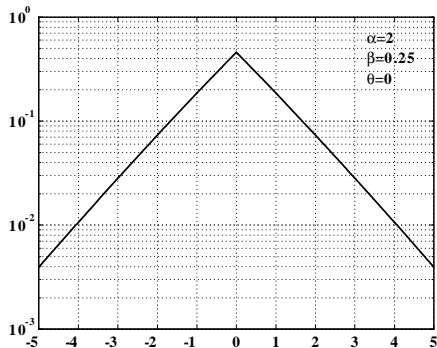
with solution

$$G_{2,\beta}^0(x, t) = \frac{1}{2} t^{-\beta/2} M_{\beta/2} \left(|x|/t^{\beta/2} \right), \quad -\infty < x < +\infty, t \geq 0$$

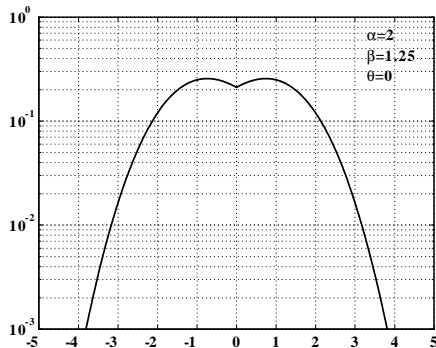
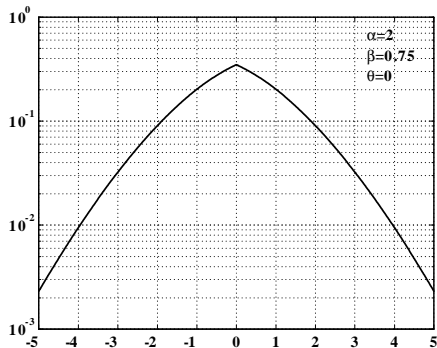
$M_{\beta/2}$ is a function of Wright type of order $\beta/2$ defined for any order $\nu \in (0, 1)$ by

$$M_\nu(z) = \sum_{n=0}^{\infty} \frac{(-z)^n}{n! \Gamma[-\nu n + (1 - \nu)]} = \frac{1}{\pi} \sum_{n=1}^{\infty} \frac{(-z)^{n-1}}{(n-1)!} \Gamma(\nu n) \sin(\pi \nu n)$$

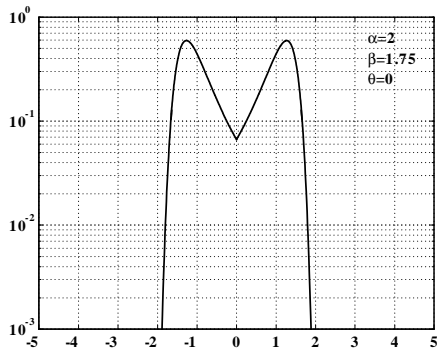
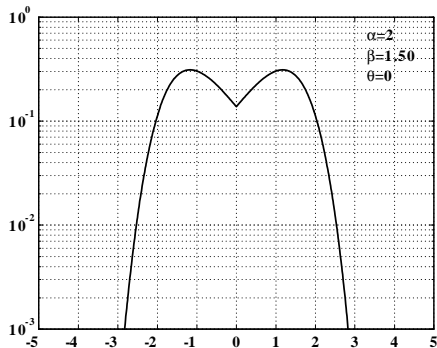
Fractional diffusion equation



Fractional diffusion equation

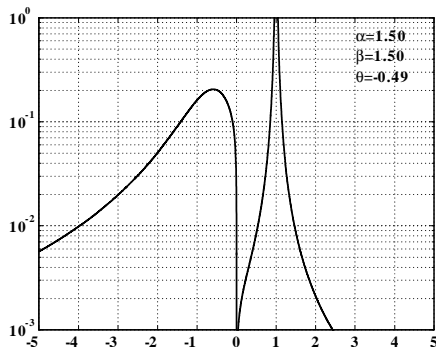
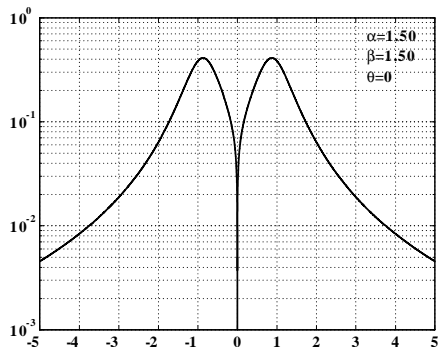


Fractional diffusion equation

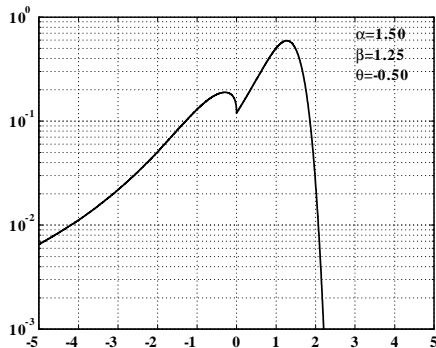
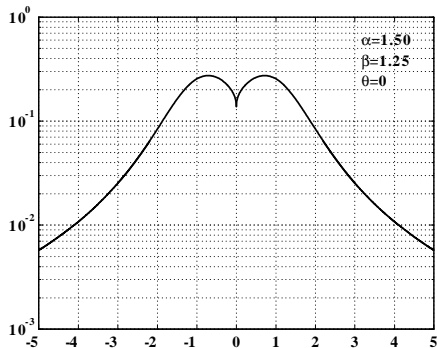


Fractional diffusion equation

Space-time fractional diffusion equation. Some examples



Fractional diffusion equation



A fractional nonlinear equation. Stochastic solution

A fractional version of the KPP equation, studied by McKean



$$\boxed{{}_t D_*^\alpha u(t, x) = \frac{1}{2} {}_x D_\theta^\beta u(t, x) + u^2(t, x) - u(t, x)}$$

${}_t D_*^\alpha$ is a Caputo derivative of order α

$${}_t D_*^\alpha f(t) = \begin{cases} \frac{1}{\Gamma(m-\beta)} \int_0^t \frac{f^{(m)}(\tau) d\tau}{(t-\tau)^{\alpha+1-m}} & m-1 < \alpha < m \\ \frac{d^m}{dt^m} f(t) & \alpha = m \end{cases}$$

${}_x D_\theta^\beta$ is a Riesz-Feller derivative defined through its Fourier symbol

$$\mathcal{F} \left\{ {}_x D_\theta^\beta f(x) \right\} (k) = -\psi_\beta^\theta(k) \mathcal{F} \{ f(x) \} (k)$$

with $\psi_\beta^\theta(k) = |k|^\beta e^{i(\text{sign } k)\theta\pi/2}$.

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$$\boxed{{}_t D_*^\alpha u(t, x) = \frac{1}{2} {}_x D_\theta^\beta u(t, x) + u^2(t, x) - u(t, x)}$$

${}_t D_*^\alpha$ is a Caputo derivative of order α

$${}_t D_*^\alpha f(t) = \begin{cases} \frac{1}{\Gamma(m-\beta)} \int_0^t \frac{f^{(m)}(\tau) d\tau}{(t-\tau)^{\alpha+1-m}} & m-1 < \alpha < m \\ \frac{d^m}{dt^m} f(t) & \alpha = m \end{cases}$$

${}_x D_\theta^\beta$ is a Riesz-Feller derivative defined through its Fourier symbol

$$\mathcal{F} \left\{ {}_x D_\theta^\beta f(x) \right\} (k) = -\psi_\beta^\theta(k) \mathcal{F} \{ f(x) \} (k)$$

with $\psi_\beta^\theta(k) = |k|^\beta e^{i(\text{sign } k)\theta\pi/2}$.

- Physically it describes a nonlinear diffusion with growing mass and in our fractional generalization it would represent the same phenomenon taking into account memory effects in time and long range correlations in space.

A fractional nonlinear equation

The first step towards a probabilistic formulation is the rewriting as an integral equation. Take the Fourier transform (\mathcal{F}) in space and the Laplace transform (\mathcal{L}) in time

$$s^\alpha \tilde{\hat{u}}(s, k) = s^{\alpha-1} \hat{u}(0^+, k) - \frac{1}{2} \psi_\beta^\theta(k) \tilde{\hat{u}}(s, k) - \tilde{\hat{u}}(s, k) + \int_0^\infty dt e^{-st} \mathcal{F}(u^2)$$

where

$$\hat{u}(t, k) = \mathcal{F}(u(t, x)) = \int_{-\infty}^{\infty} e^{ikx} u(t, x)$$

$$\tilde{\hat{u}}(s, k) = \mathcal{L}(\hat{u}(t, k)) = \int_0^\infty e^{-st} \hat{u}(t, k)$$

This equation holds for $0 < \alpha \leq 1$ or for $0 < \alpha \leq 2$ with $\frac{\partial}{\partial t} u(0^+, x) = 0$.

Solving for $\tilde{\hat{u}}(s, k)$ one obtains an integral equation

$$\tilde{\hat{u}}(s, k) = \frac{s^{\alpha-1}}{s^\alpha + \frac{1}{2} \psi_\beta^\theta(k)} \hat{u}(0^+, k) + \int_0^\infty dt \frac{e^{-st}}{s^\alpha + \frac{1}{2} \psi_\beta^\theta(k)} \mathcal{F}(u^2(t, x))$$

A fractional nonlinear equation

Taking the inverse Fourier and Laplace transforms

$$\begin{aligned} & u(t, x) \\ = & \boxed{E_{\alpha,1}(-t^\alpha)} \int_{-\infty}^{\infty} dy \mathcal{F}^{-1} \left(\frac{E_{\alpha,1} \left(- \left(1 + \frac{1}{2} \psi_\beta^\theta(k) \right) t^\alpha \right)}{E_{\alpha,1}(-t^\alpha)} \right) (x-y) u(0, y) \\ & + \int_0^t d\tau \boxed{(t-\tau)^{\alpha-1} E_{\alpha,\alpha}(-(t-\tau)^\alpha)} \\ & \int_{-\infty}^{\infty} dy \mathcal{F}^{-1} \left(\frac{E_{\alpha,\alpha} \left(- \left(1 + \frac{1}{2} \psi_\beta^\theta(k) \right) (t-\tau)^\alpha \right)}{E_{\alpha,\alpha}(-(t-\tau)^\alpha)} \right) (x-y) u^2(\tau, y) \end{aligned}$$

$E_{\alpha,\rho}$ is the generalized Mittag-Leffler function $E_{\alpha,\rho}(z) = \sum_{j=0}^{\infty} \frac{z^j}{\Gamma(\alpha j + \rho)}$

$$E_{\alpha,1}(-t^\alpha) + \int_0^t d\tau (t-\tau)^{\alpha-1} E_{\alpha,\alpha}(-(t-\tau)^\alpha) = 1$$

A fractional nonlinear equation

We define the following propagation kernel

$$G_{\alpha,\rho}^{\beta}(t,x) = \mathcal{F}^{-1} \left(\frac{E_{\alpha,\rho} \left(- \left(1 + \frac{1}{2} \psi_{\beta}^{\theta}(k) \right) t^{\alpha} \right)}{E_{\alpha,\rho}(-t^{\alpha})} \right) (x)$$

$$\begin{aligned} & u(t,x) \\ = & \boxed{E_{\alpha,1}(-t^{\alpha})} \int_{-\infty}^{\infty} dy \boxed{G_{\alpha,1}^{\beta}(t,x-y)} u(0^{+},y) \\ & + \int_0^t d\tau \boxed{(t-\tau)^{\alpha-1} E_{\alpha,\alpha}(-(t-\tau)^{\alpha})} \\ & \int_{-\infty}^{\infty} dy \boxed{G_{\alpha,\alpha}^{\beta}(t-\tau,x-y)} u(\tau,y) \end{aligned}$$

$E_{\alpha,1}(-t^{\alpha})$ and $(t-\tau)^{\alpha-1} E_{\alpha,\alpha}(-(t-\tau)^{\alpha})$ = survival probability up to time t and the probability density for the branching at time τ (branching process B_{α})

A fractional nonlinear equation

The propagation kernels satisfy the conditions to be the Green's functions of stochastic processes in \mathbb{R} :

$$u(t, x) = \mathbb{E}_x (u(0^+, x + \xi_1) u(0^+, x + \xi_2) \cdots u(0^+, x + \xi_n))$$

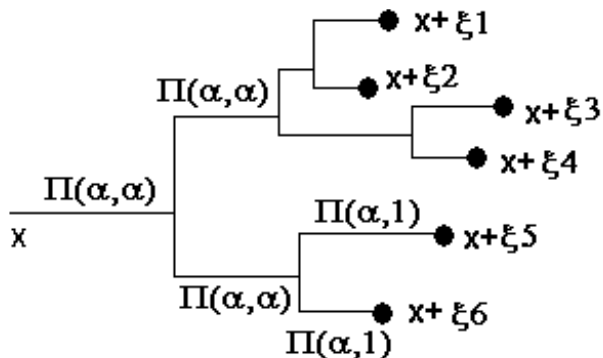
Denote the processes associated to $G_{\alpha,1}^\beta(t, x)$ and $G_{\alpha,\alpha}^\beta(t, x)$, respectively by $\Pi_{\alpha,1}^\beta$ and $\Pi_{\alpha,\alpha}^\beta$

Theorem: *The nonlinear fractional partial differential equation, with $0 < \alpha \leq 1$, has a stochastic solution, the coordinates $x + \xi_i$ in the arguments of the initial condition obtained from the exit values of a propagation and branching process, the branching being ruled by the process B_α and the propagation by $\Pi_{\alpha,1}^\beta$ for the first particle and by $\Pi_{\alpha,\alpha}^\beta$ for all the remaining ones.*

A sufficient condition for the existence of the solution is

$$|u(0^+, x)| \leq 1$$

A fractional nonlinear equation



Geometric interpretation of fractional integration: shadows on the walls

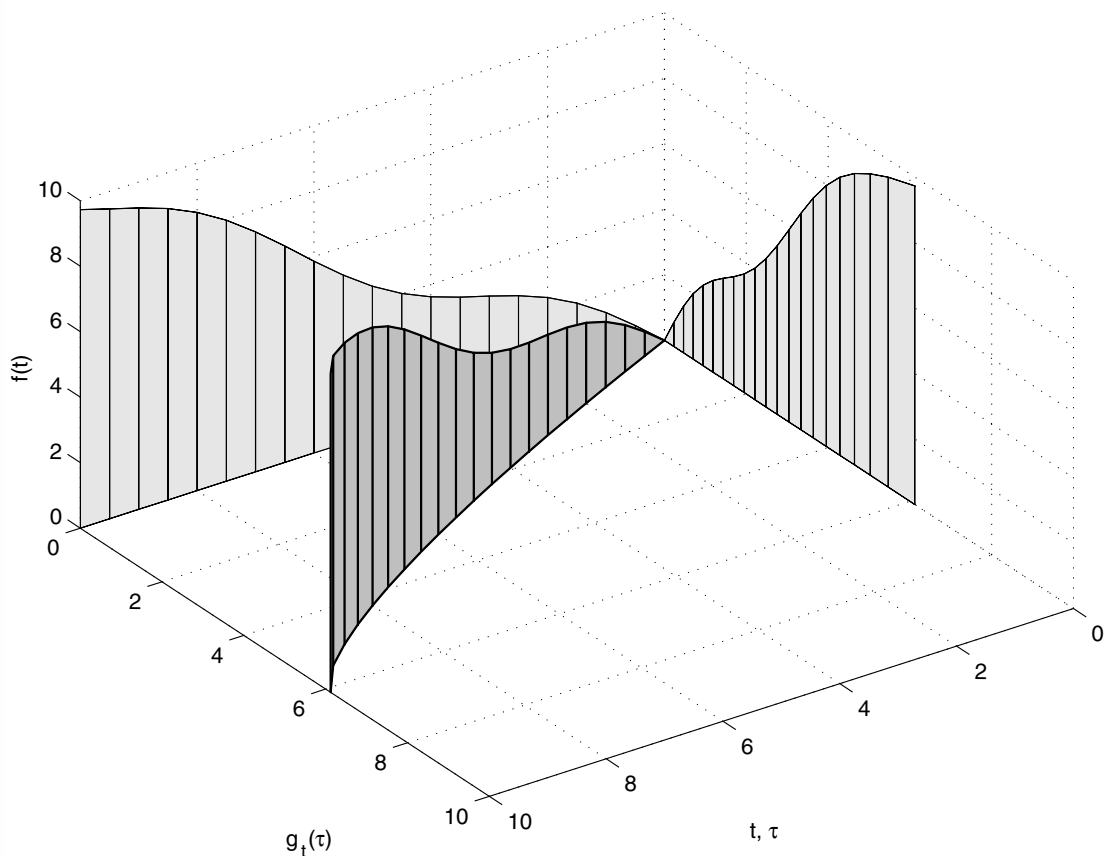
$${}_0I_t^\alpha f(t) = \frac{1}{\Gamma(\alpha)} \int_0^t f(\tau)(t - \tau)^{\alpha-1} d\tau, \quad t \geq 0,$$

$${}_0I_t^\alpha f(t) = \int_0^t f(\tau) dg_t(\tau),$$

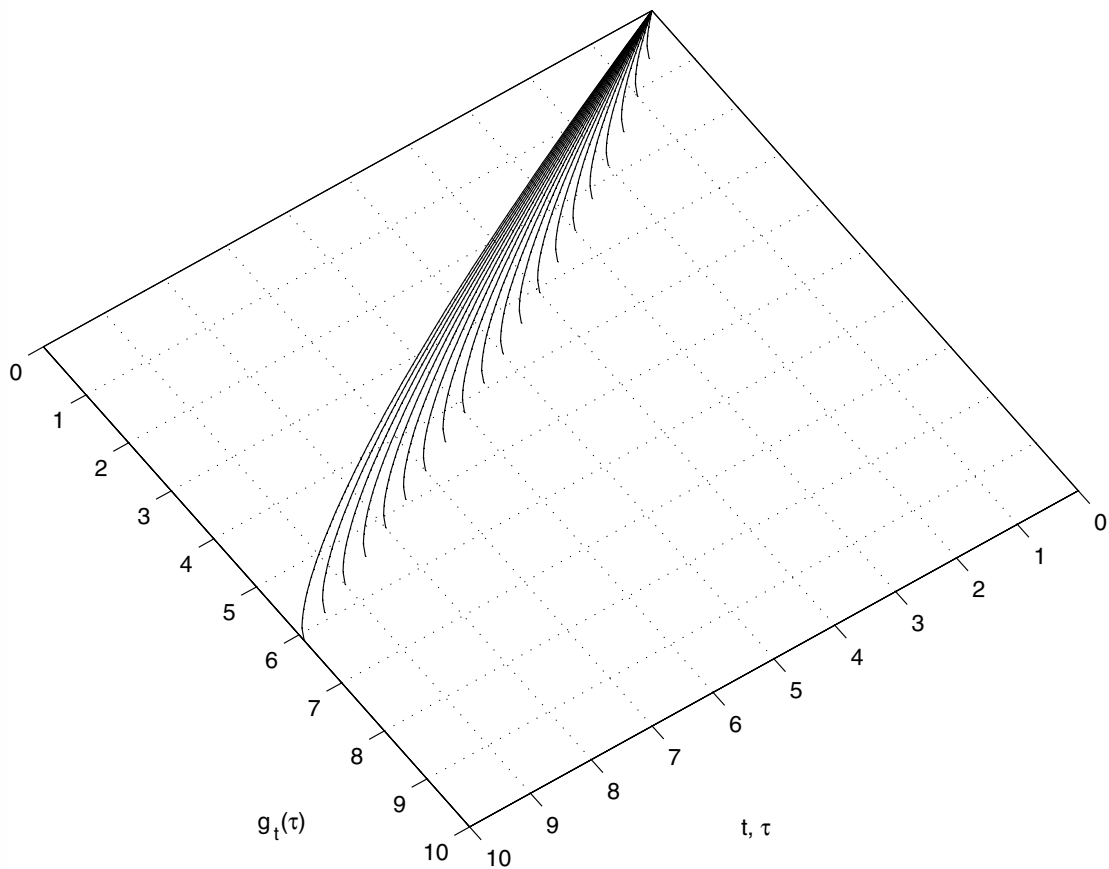
$$g_t(\tau) = \frac{1}{\Gamma(\alpha + 1)} \{t^\alpha - (t - \tau)^\alpha\}.$$

For $t_1 = kt$, $\tau_1 = k\tau$ ($k > 0$) we have:

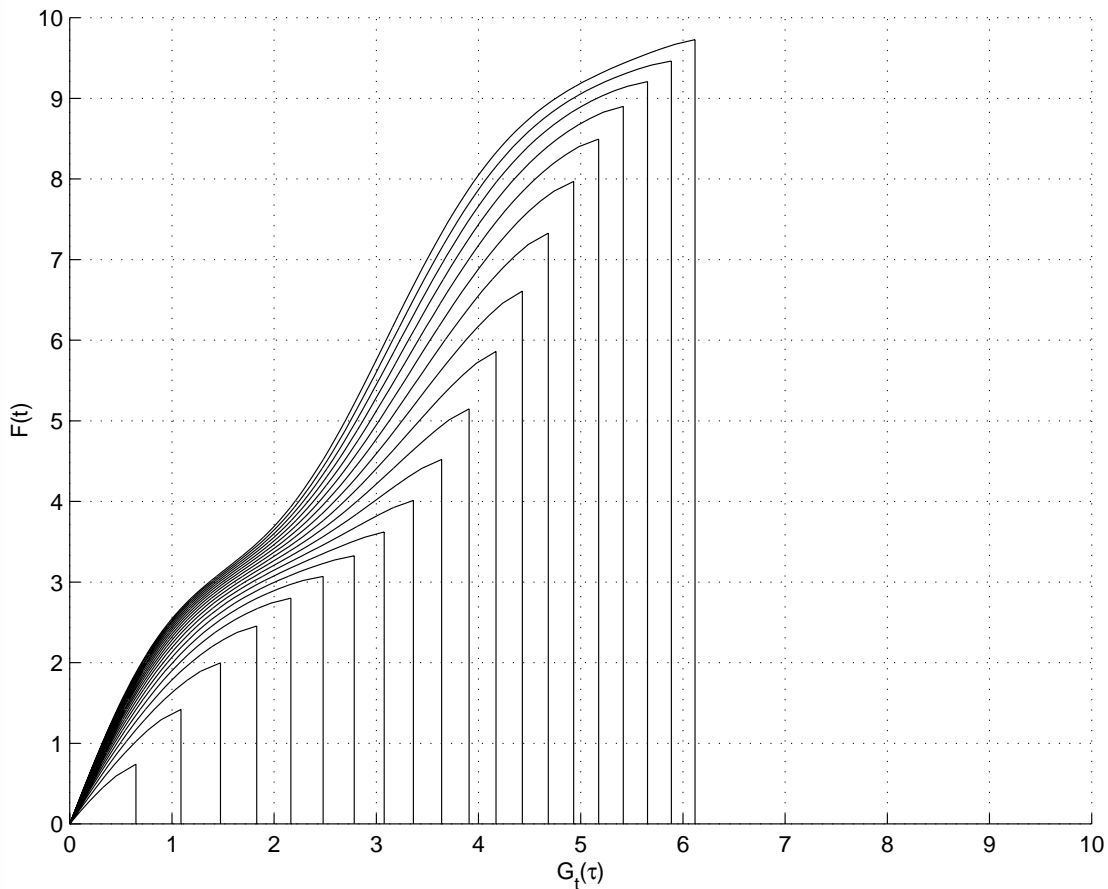
$$g_{t_1}(\tau_1) = g_{kt}(k\tau) = k^\alpha g_t(\tau).$$



“Live fence” and its shadows: ${}_0I_t^1 f(t)$ and ${}_0I_t^\alpha f(t)$,
for $\alpha = 0.75$, $f(t) = t + 0.5 \sin(t)$, $0 \leq t \leq 10$.



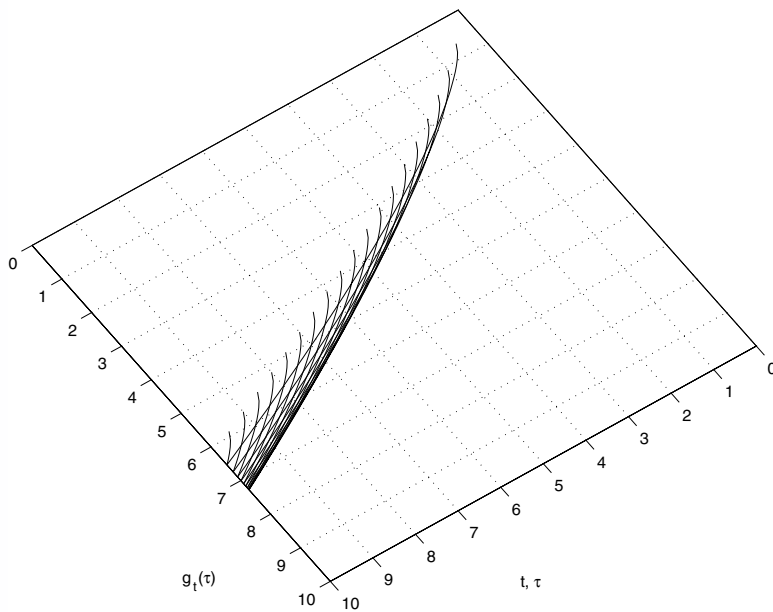
“Live fence”: basis shape is changing
for ${}_0I_t^\alpha f(t)$, $\alpha = 0.75$, $0 \leq t \leq 10$.



Snapshots of the changing “shadow” of changing “fence” for ${}_0I_t^\alpha f(t)$, $\alpha = 0.75$, $f(t) = t + 0.5 \sin(t)$, with the time interval $\Delta t = 0.5$ between the snapshots.

Right-sided R-L integral

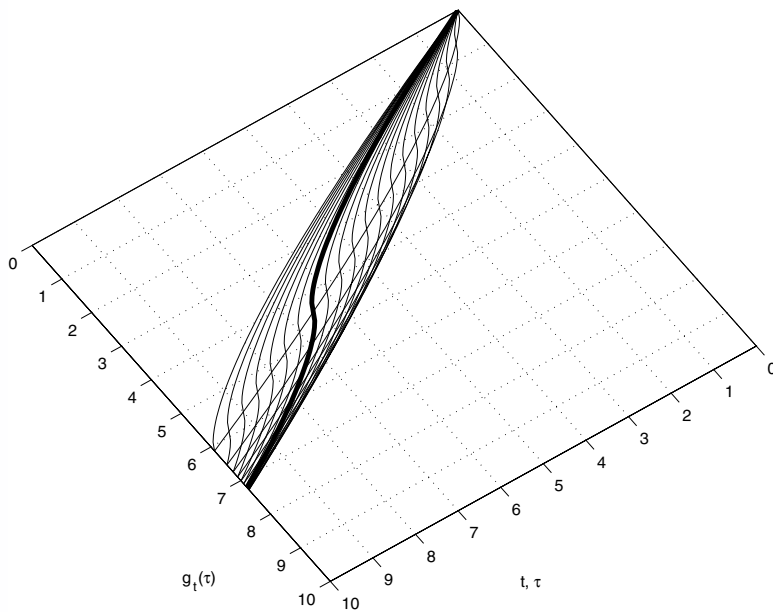
$${}_tI_0^\alpha f(t) = \frac{1}{\Gamma(\alpha)} \int_t^b f(\tau)(\tau - t)^{\alpha-1} d\tau, \quad t \leq b,$$



$${}_tI_{10}^\alpha f(t), \quad \alpha = 0.75, \quad 0 \leq t \leq 10$$

Riesz potential

$${}_0R_b^\alpha f(t) = \frac{1}{\Gamma(\alpha)} \int_0^b f(\tau) |\tau - t|^{\alpha-1} d\tau, \quad 0 \leq t \leq b,$$



$${}_0R_{10}^\alpha f(t), \quad \alpha = 0.75, \quad 0 \leq t \leq 10$$

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